High Density QCD

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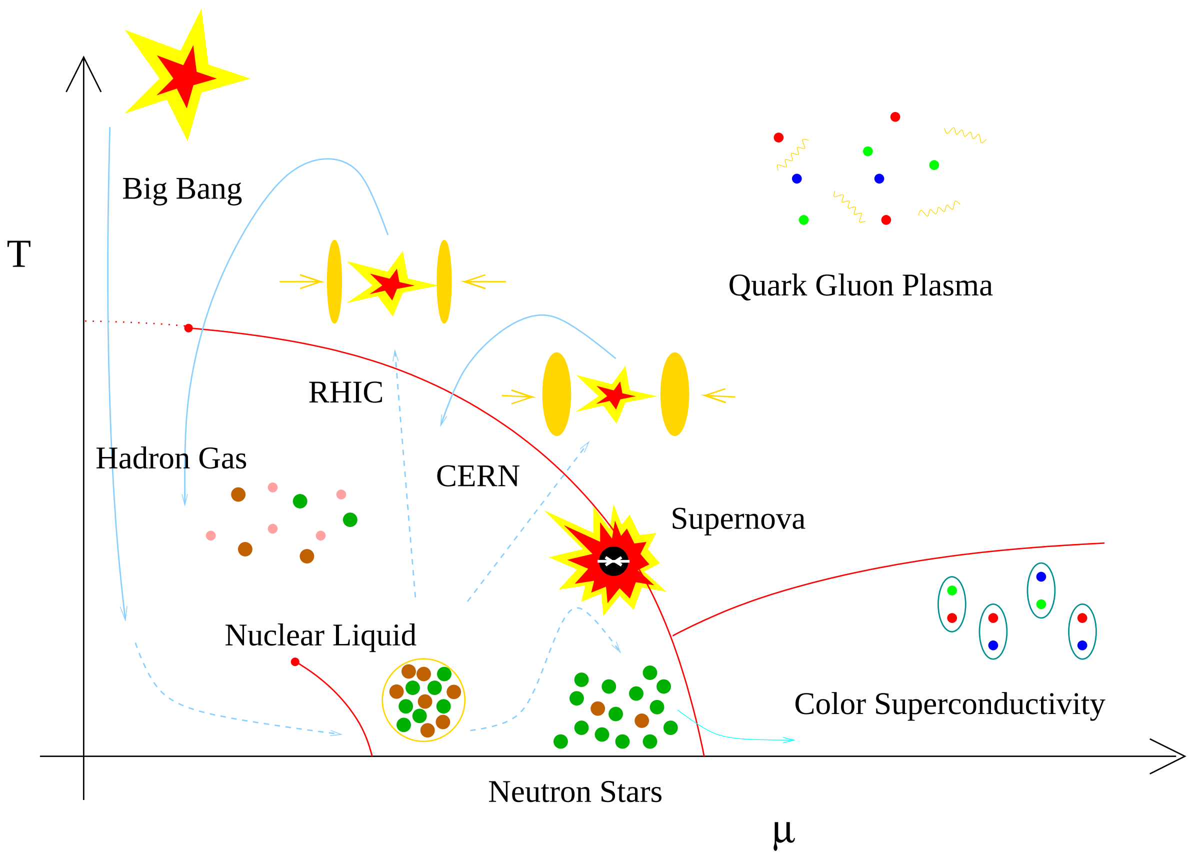
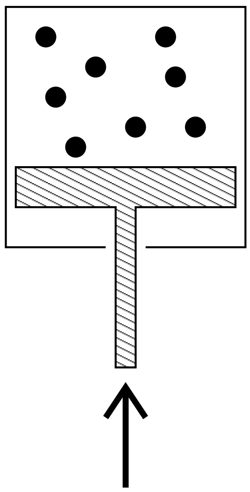
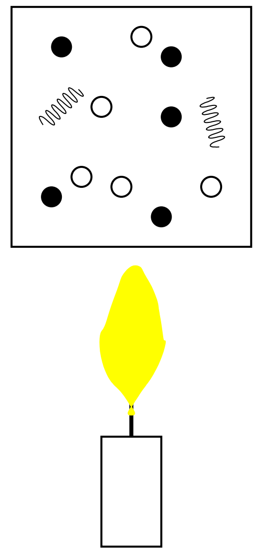


QCD in Xtreme Conditions

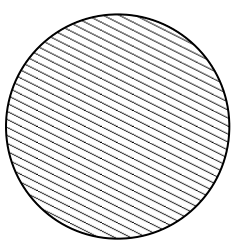
High Density



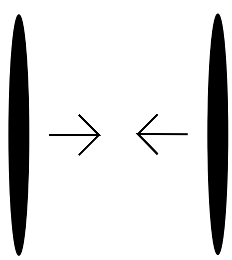
High Temperature



more practical:



neutron
stars



relativistic
heavy−ion
collisions

big bang

WHY?

- These conditions exist in the universe:

  Neutron Stars, Big Bang

- Exploring the entire phase diagram is important to understanding the phase that we happen to live in

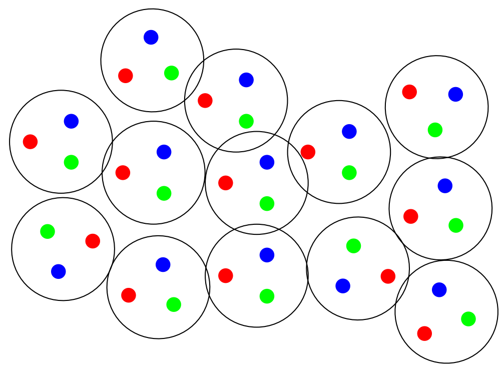
  We cannot understand the structure of hadrons without understanding the non-trivial QCD vacuum.
  We cannot understand the vacuum state without understanding how it can be modified.

- QCD simplifies in extreme environments:

  Study QCD matter in a regime where quarks and gluons are indeed the correct degrees of freedom.

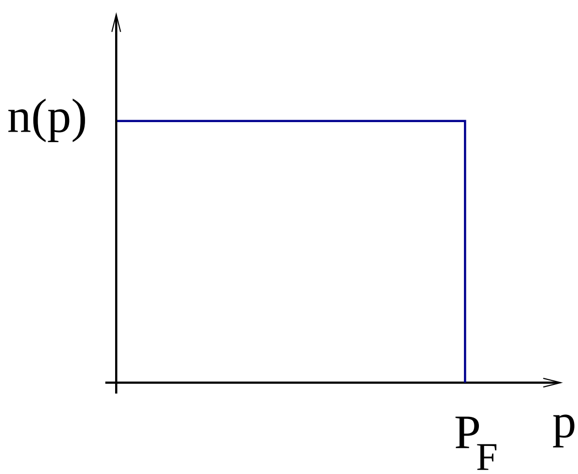
Dense Hadronic Matter

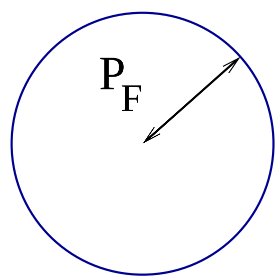
- consider baryon density $n_B \gg 1\,\mathrm{fm}^{-3}$



quarks expected to move freely

- groundstate: cold quark matter (quark fermi liquid)

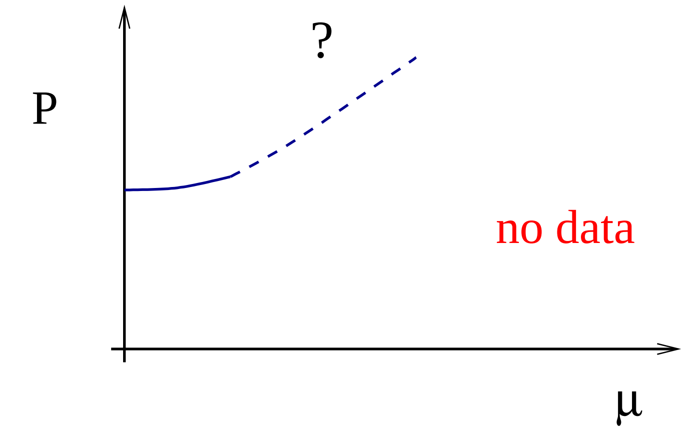


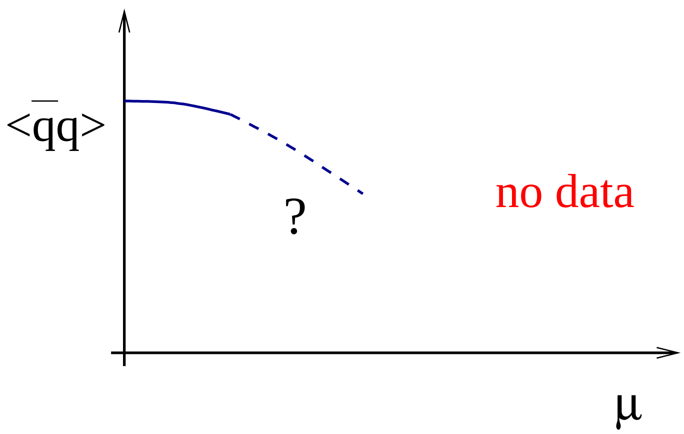


- only quarks with $p \sim p_F \gg \Lambda_{QCD}$ scatter
- $p_F \gg \Lambda_{QCD} \to$ effective coupling is weak

$\to$ no chiral symmetry breaking
$\to$ no confinement
$\to$ no dynamically generated masses

Lattice Results



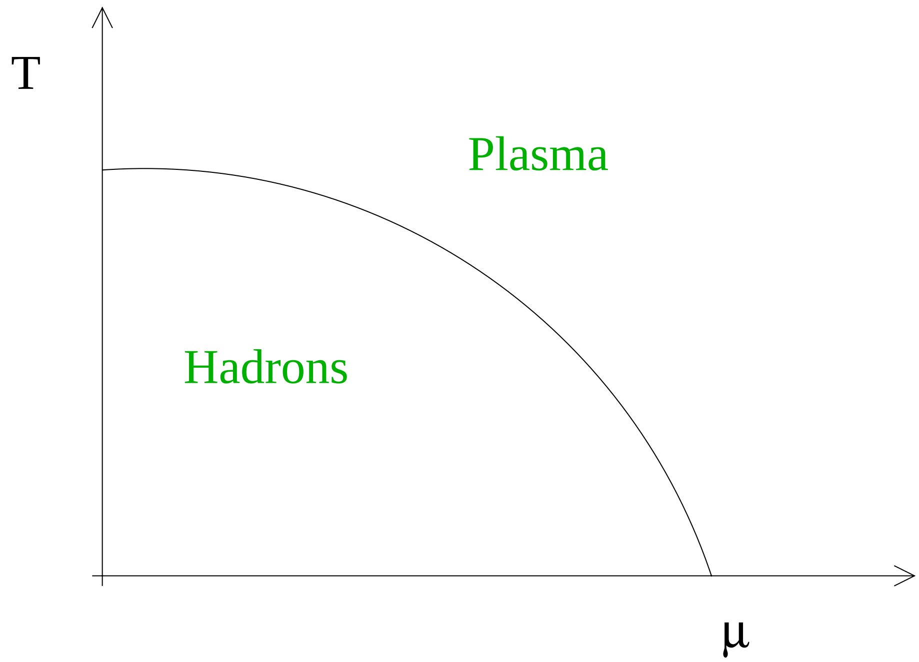


- Problem: No lattice data

$$Z = \int dA \underbrace{\det(\not{D} + i\mu\gamma_4)}_{\text{complex}} e^{-S}$$

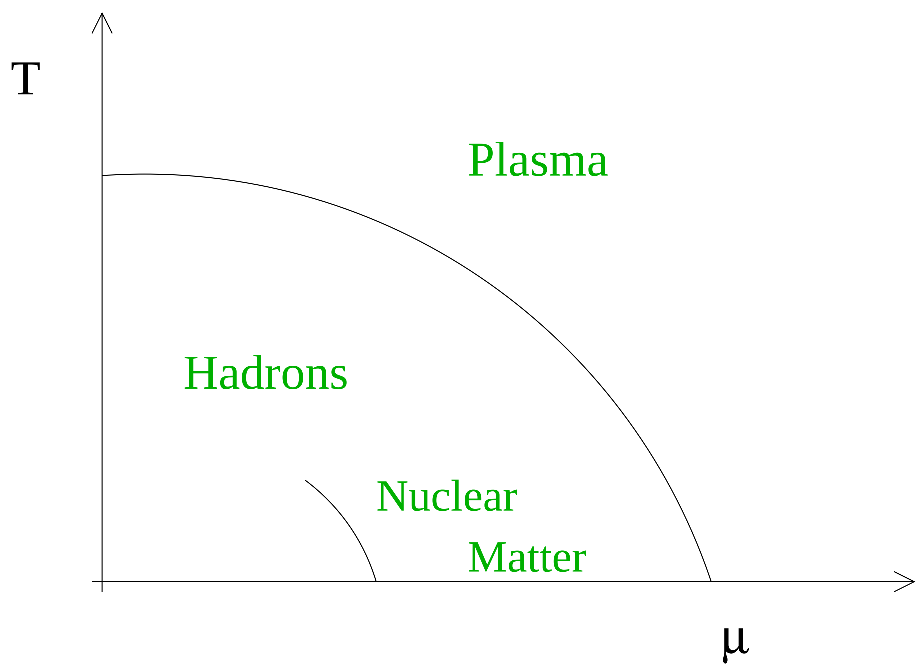
- "sign" problem: no importance sampling
- progress in simulating small μ, $T \sim T_c$
  - improved reweighting (Fodor, Katz)
  - imaginary chemical potential (Forcrand, Philipsen, ...)
  - Taylor expansion (Karsch et al.)

Phase Diagram: First Try



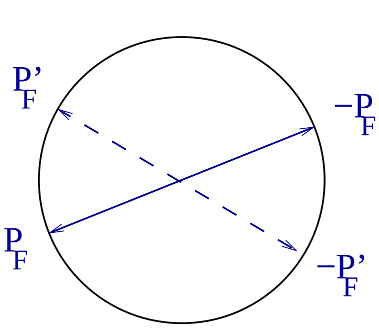
Where are we?

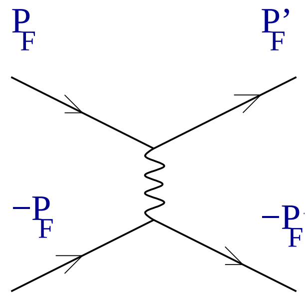
Phase Diagram: Second Try



Superfluidity

- dominant interaction





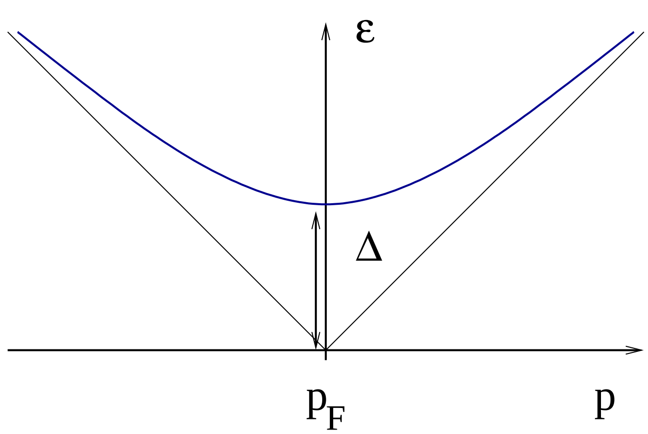
Uses the whole Fermi surface coherently

- very degenerate perturbation theory

$\Rightarrow$ instability (BCS) in weak coupling

$\Rightarrow$ $q(p_F)q(-p_F)$ condensate

- arbitrarily weak (attractive!) interaction leads to superfluidity/superconductivity (pairs neutral/charged)



$$\epsilon(p) = \sqrt{(p - p_F)^2 + \Delta^2}$$

$\rightarrow$ gap in spectrum, transport without dissipation, ...

- QCD: gluon exchange attractive in $\bar{3}$ channel

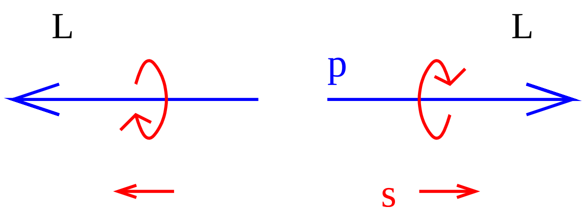
$$3 \times 3 = 6_S + 3_A$$ Flux reduced $\Rightarrow$ attractive

- spin-flavor-color wave function

$$(\uparrow\downarrow - \downarrow\uparrow) \times (ud - du) \times (rb - br) \qquad s = 0,\ I = 0,\ c = \bar{3}$$

order parameter: $\Phi^a = \epsilon^{abc}\langle q^b C\gamma_5 \tau_2 q^c\rangle$

- chiral symmetry is not broken



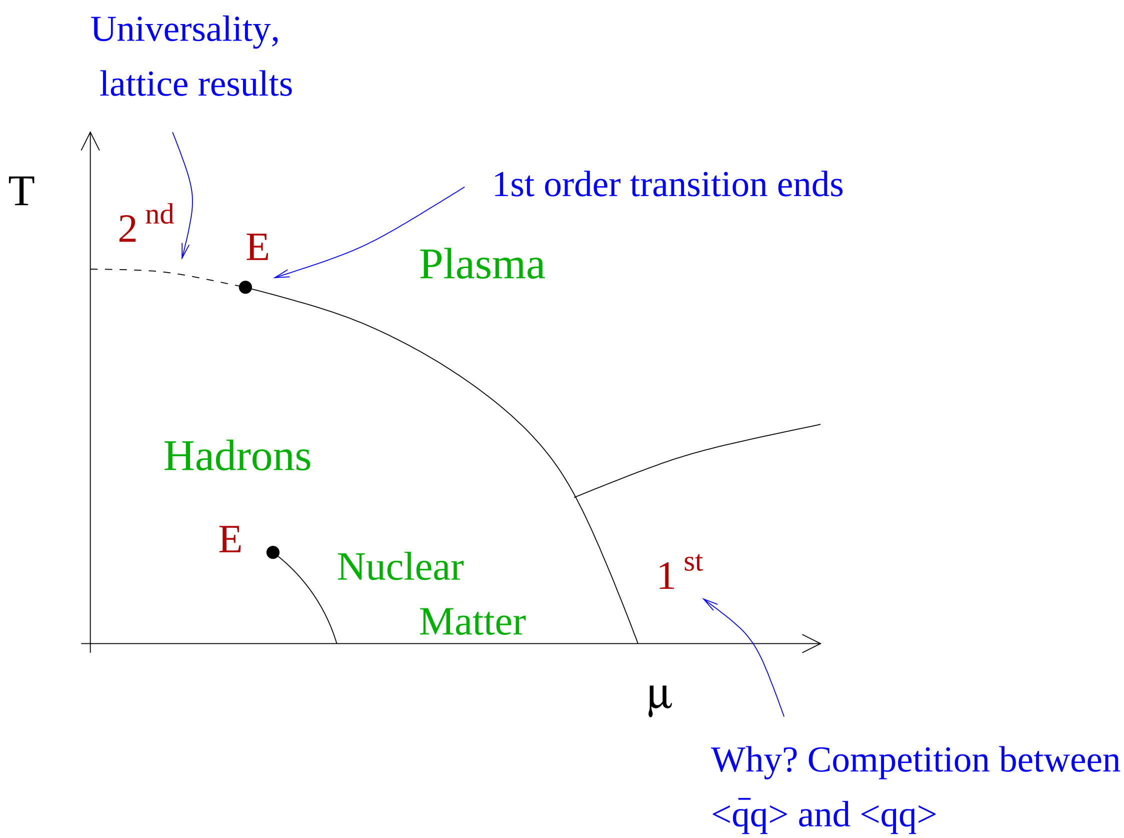
$$\Phi \sim \langle q_L q_L\rangle - \langle q_R q_R\rangle$$

- color symmetry broken by Higgs mechanism (Meissner effect)

$$\Phi^a \in [\bar{3}];\ \Rightarrow\ SU(3) \to SU(2)$$

  – 5/8 gluons acquire mass via Higgs mechanism
  – $SU(2)$ is confined

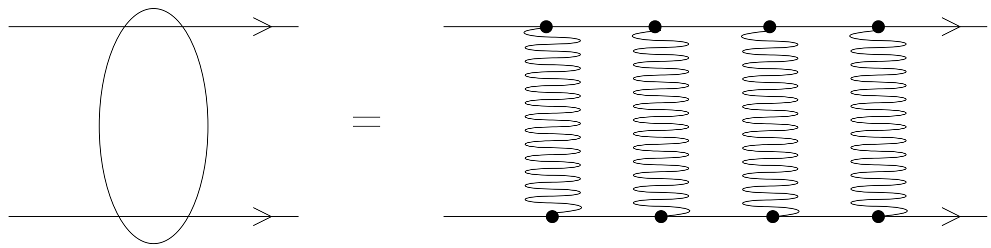
Phase Diagram: Second Revision



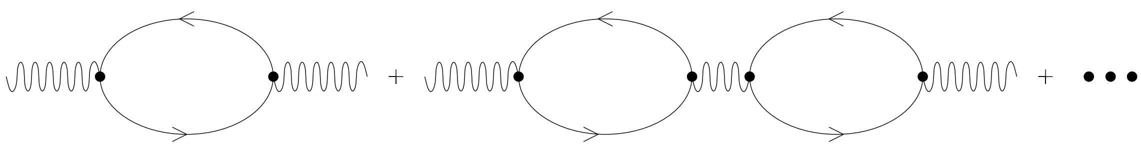
- critical endpoint (E) persists even if $m \neq 0$
- Critical fluctuations observable in heavy ion collisions?
  - charge fluctuations
  - baryon number fluctuations
  - p_T fluctuations, ...

Gap Equation

- $\mu \gg \Lambda_{QCD}$: perturbative forces dominate



- gluon exchange: small angle scattering dominates $\rightarrow$ medium effects important



$$D_E = \frac{1}{\vec{q}^2 + 2m^2} \qquad m^2 = \frac{N_f}{4\pi^2} g^2 \mu^2 \qquad \text{Debye screening}$$

$$D_M = \frac{1}{\vec{q}^2 + i\frac{\pi}{2} m \frac{\omega}{q}} \qquad \omega < q \qquad \text{Landau damping}$$

- consider $\omega \simeq \Delta$. Typical momenta

$$q_E \simeq g\mu, \qquad q_M \simeq (g^2\mu^2\Delta)^{1/3}$$

Superconductivity driven by magnetic forces!

- Note: Perturbation theory is self-consistent

$$\mu \rightarrow \infty \quad \Rightarrow \quad q_M \gg \Lambda_{QCD}$$

Eliashberg Equation

- retardation important: Eliashberg theory

$$\Delta(p_0) = \frac{g^2}{18\pi^2}\int dq_0 \left\{\log\left(\frac{b_M}{|p_0-q_0|}\right) + \ldots\right\} \frac{\Delta(q_0)}{\sqrt{q_0^2+\Delta(q_0)^2}}$$

colinear log BCS log

- double logarithmic behavior

$$\Delta_0 \sim \frac{g^2}{18\pi^2}\Delta_0 \left[\log\left(\frac{\mu}{\Delta_0}\right)\right]^2$$

$$\rightarrow \Delta_0 \sim \exp\left(-\frac{c}{\sqrt{\alpha_s}}\right)$$

(BCS result $\Delta_0 \sim \exp\left(-\frac{c}{\alpha_s}\right)$)

- more careful analysis

$$c = 3\pi^2/\sqrt{2}$$

- condensation energy

$$\epsilon \simeq N\Delta_0^2\left(\frac{\mu^2}{4\pi^2}\right)$$

- typical numbers

$$\Delta_0 \simeq 100 \ \text{MeV} \ (\rho \sim 5\rho_0)$$

$$T_c \simeq 50 \ \text{MeV}$$

Beautiful Case: $N_f = 3$

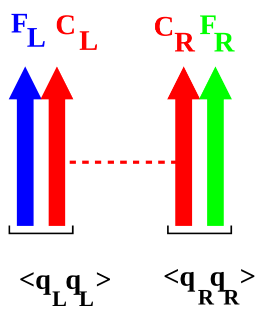
- quarks q_i^a $[i = (u, d, s);\ a = (r, b, g)]$. Order parameter?

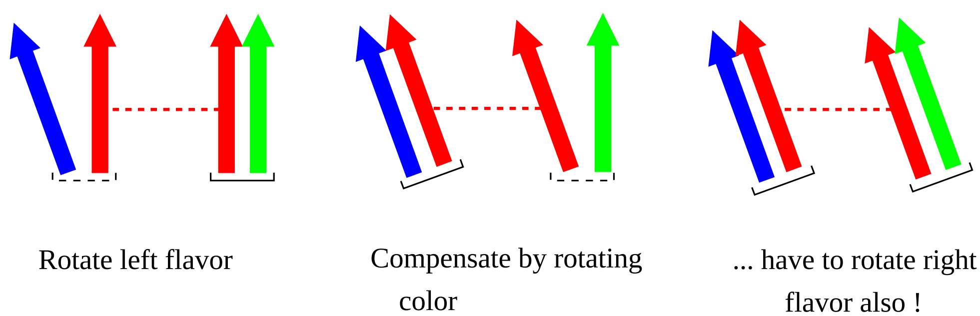
$$\langle q_i^a q_j^b \rangle = \phi\, \epsilon^{abI} \epsilon_{ijI} \quad \text{Color} - \text{Flavor} - \text{Locking}$$

$$\begin{array}{ll} \langle ud \rangle = \langle us \rangle = \langle ds \rangle & \text{Color and Flavor} \\ \langle rb \rangle = \langle rg \rangle = \langle bg \rangle & \text{direction aligned} \end{array}$$

- Symmetry breaking pattern:

$$SU(3)_L \times SU(3)_R \times [SU(3)]_C \times U(1) \to SU(3)_{C+F}$$





- Novel mechanism for Chiral Symmetry Breaking:

$$\langle \psi_L \psi_L \rangle = -\langle \psi_R \psi_R \rangle$$

Breaks chiral $SU(3)_L \times SU(3)_R$ symmetry because

$$SU(3)_L \overset{Lock}{\longleftrightarrow} SU(3)_C \overset{Lock}{\longleftrightarrow} SU(3)_R$$

- All quarks and gluons acquire a mass gap

CFL Phase: Excitations

- quark pairs are charged, but $U(1)_{EM}$ remains unbroken

$$\gamma^* = \alpha_{11}\gamma + \alpha_{12}g_3 + \alpha_{13}g_8$$

Material is a transparent insulator

- excitations

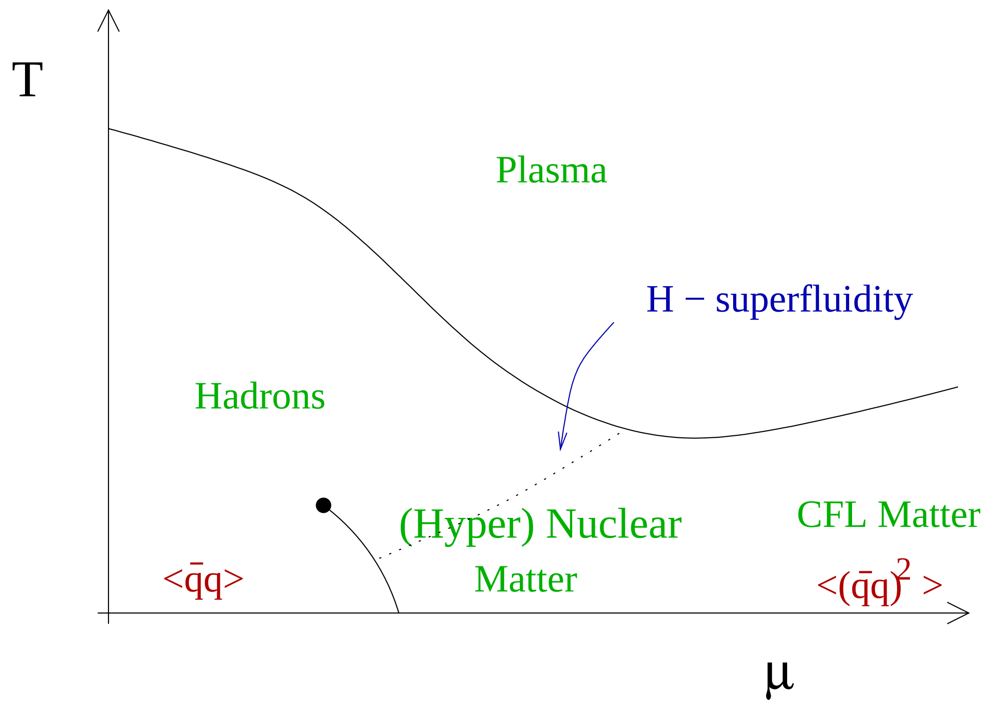
[8]	$Q = 0, \pm 1$	Goldstones (π, K, η)	$(QQ)(QQ)^{-1}, \ldots$
[8] + [1]	$Q = 0, \pm 1$	baryons (p, n, Λ, Σ)	$(Q)(QQ), \ldots$
[8]	$Q = 0, \pm 1$	vectors $(\rho, K^*, \ldots)$	$g, QQ^{-1}, \ldots$
[1]	$Q = 0$	$U(1)_A$ Goldstone (η')	$QQ^{-1}, (QQ)(QQ)^{-1}, \ldots$
[1]	$Q = 0$	$U(1)_B$ Goldstone	$QQ^{-1}, (QQ)(QQ)^{-1}, \ldots$

- order parameters

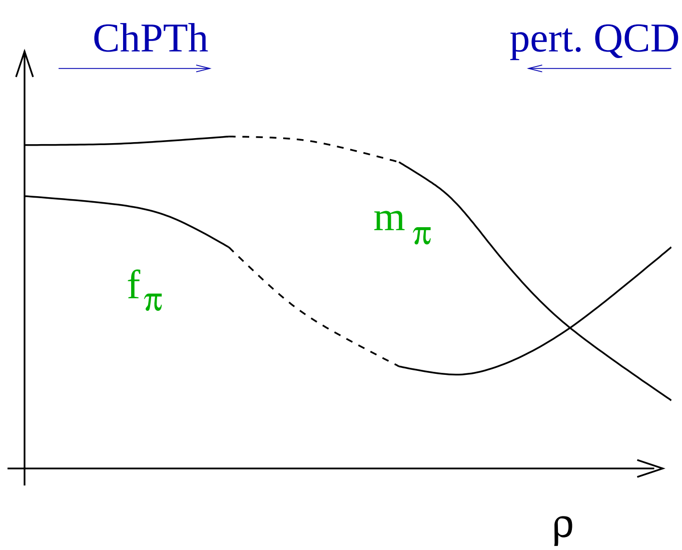
$\langle \bar{q}q \rangle$	chiral symmetry breaking
$\langle (uds)(uds) \rangle$	H superfluidity

Continuity between quark and hadron matter?

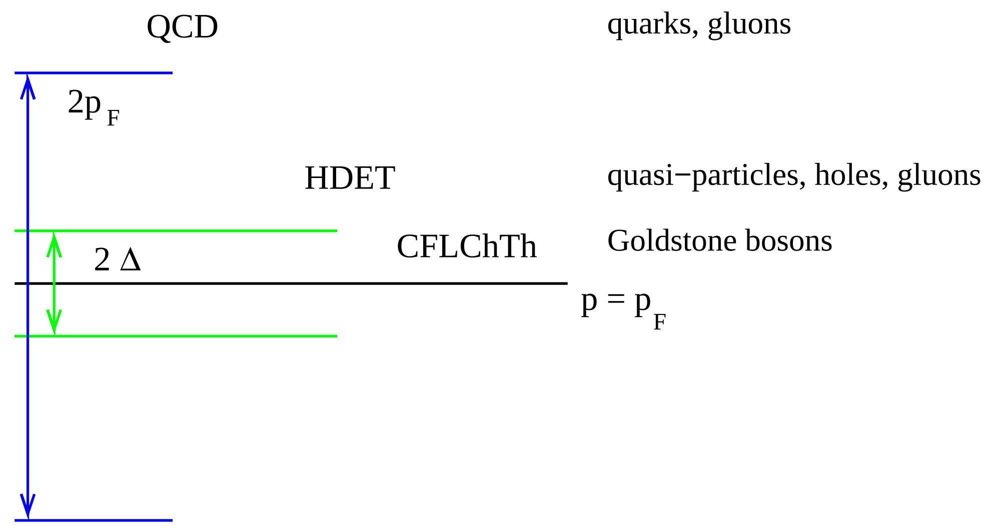
Phase Diagram: Third Revision



Density dependence of hadronic parameters calculabale at both low density (effective theory) and high density (perturbative QCD).



Effective Field Theories

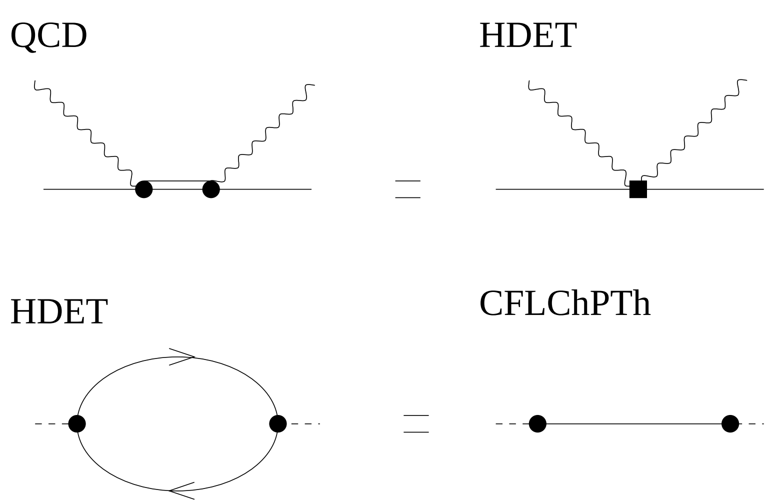


QCD $\mathcal{L} = \bar\psi(i\not{D} + \mu\gamma_0)\psi - \frac{1}{4}G^a_{\mu\nu}G^a_{\mu\nu}$

HDET $\mathcal{L} = \psi_v^\dagger(iv\cdot D)\psi_v - \frac{\Delta}{2}\psi_{-v}^T C\psi_v - \frac{1}{4}G^a_{\mu\nu}G^a_{\mu\nu} + \ldots$

CFLChPTh $\mathcal{L} = \frac{f_\pi^2}{4}\mathrm{Tr}\left[\nabla_0\Sigma\nabla_0\Sigma^\dagger - v^2\vec\nabla\Sigma\vec\nabla\Sigma^\dagger\right] + \ldots$

- coefficients determined by matching Green functions

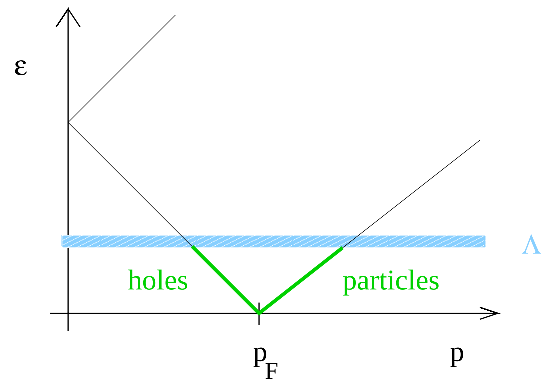


- Applications: – $f_\pi, m_\pi, \ldots$ (collective modes)
  - – phase structure as a function of $m_s, \mu_e, \ldots$
  - – transport properties, ...

High density effective theory

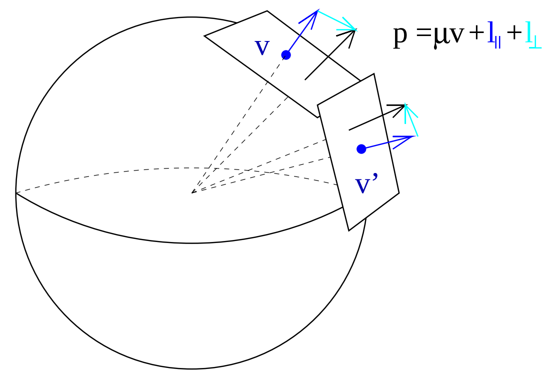
- quasi-particles (holes)

$$E_\pm = -\mu + \sqrt{\vec{p}^2 + m^2}$$
$$\simeq -\mu \pm |\vec{p}|$$



- effective field theory on v-patches

$$\psi_{v\pm} = e^{-i\mu v\cdot x}\left(\frac{1 \pm \vec{\alpha}\cdot\vec{v}}{2}\right)\psi$$

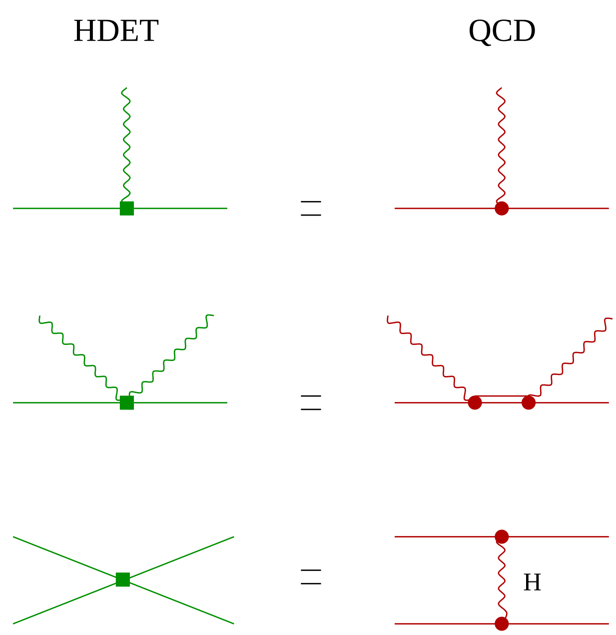


- effective lagrangian for ψ_{v+}

$$\mathcal{L} = \psi_v^\dagger (iv \cdot D)\psi_v$$

$$+ \frac{1}{2p_F}\psi_v^\dagger \left[(\vec{\alpha}_\perp \cdot \vec{D})^2 + MM^\dagger\right]\psi_v$$

$$+ \frac{\Gamma_{\vec{v}\cdot\vec{v}'}}{p_F^2}(\psi_{v'}^\dagger \psi_v)(\psi_{v'}^\dagger \psi_v) + \ldots$$

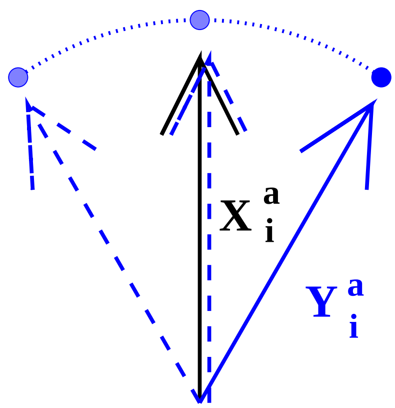


- superconducting phase:

$$\mathcal{L} \rightarrow \mathcal{L} + \psi_v C \Delta \psi_{-v} + \text{h.c.}$$

Effective Chiral Theory

- How to set up a color-flavor wave:



$$X_i^a = \epsilon_{ijk}\epsilon^{abc}\langle(\psi_L)_j^b C(\psi_L)_k^c\rangle$$

$$Y_i^a = \epsilon_{ijk}\epsilon^{abc}\langle(\psi_R)_j^b C(\psi_R)_k^c\rangle$$

$$\langle X_i^a\rangle \sim \delta_i^a$$

$$\langle Y_i^a\rangle \sim \delta_i^a$$

- low energy degrees of freedom

$$\Sigma = XY^\dagger = \exp\left(\frac{i\phi^a\lambda^a}{f_\pi}\right) \qquad \phi^a = (\pi, K, \eta, \eta')$$

e.g. $K^0 \sim \epsilon^{abc}\epsilon_{ade}(\bar{u}_R^b C \bar{s}_R^c)(d_L^d C u_L^e)$

- effective theory

$$\begin{aligned}\mathcal{L} \;=\; & \tfrac{f_\pi^2}{4}\mathrm{Tr}\left(\nabla_0\Sigma\nabla_0\Sigma^\dagger - v^2\vec{\nabla}\Sigma\vec{\nabla}\Sigma^\dagger\right)\\ & + A\mathrm{Tr}(M\Sigma^\dagger)Ve^{i\Theta} \qquad \leftarrow \quad U_A(1) \text{ anomaly}\\ & + B_1\left[\mathrm{Tr}(M\Sigma^\dagger)\right]^2 + \ldots\end{aligned}$$

- expansion

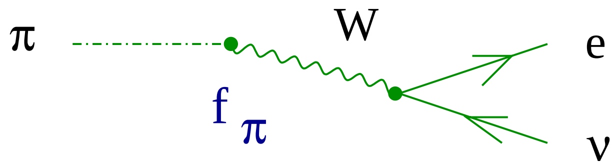
$$\mathcal{L} \sim f_\pi^2\Delta^2\left(\frac{\vec{\partial}}{\Delta}\right)^{N_1}\left(\frac{\partial_0 + MM^\dagger/p_F}{\Delta}\right)^{N_2}\left(\frac{MM}{p_F^2}\right)^{N_3}(\Sigma)^{N_4}(\Sigma^\dagger)^{N_4}$$

Note: $\Lambda_{\chi SB} \sim \Delta \ll 4\pi f_\pi$

Matching, part I

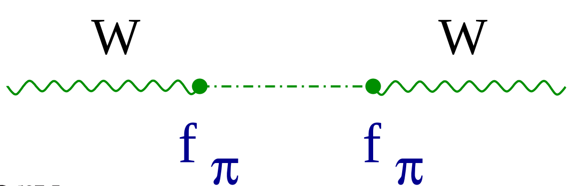
- Compute f_π: Gauge $SU(3)_L \times SU(3)_R$ flavor symmetry

$$\partial_\mu \Sigma \to \nabla_\mu \Sigma = \partial_\mu \Sigma - i W_\mu^L \Sigma + i \Sigma W_\mu^R$$

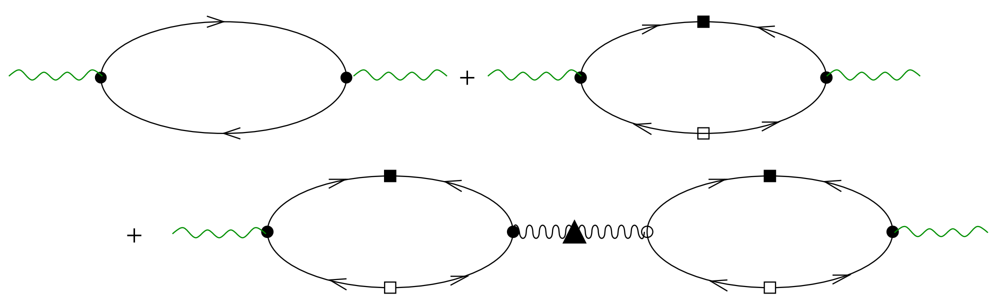


- Higgs phenomenon

$$\mathcal{L} = \frac{f_\pi^2}{4} \mathrm{Tr}\left[(W_0^L - W_0^R)^2 \right] + \ldots \qquad m_W^2 = f_\pi^2$$



- microscopic theory



$$f_\pi^2 = \frac{21 - 8\log(2)}{18} \left(\frac{\mu^2}{2\pi^2} \right)$$

- Note: $4\pi f_\pi \sim p_F \gg \Delta$

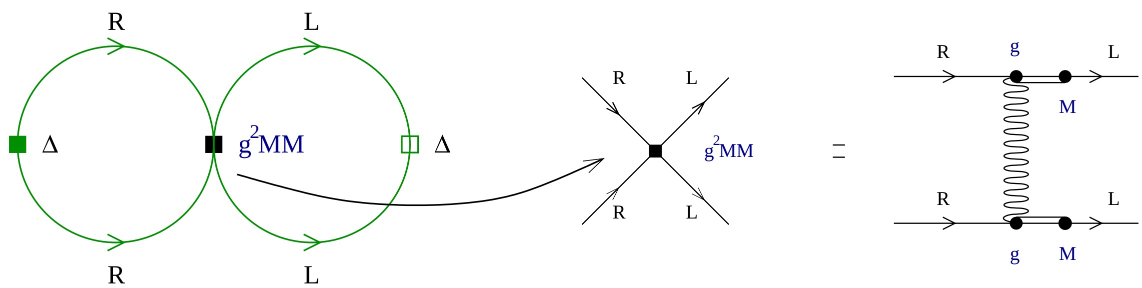
$$\rho \sim 5\rho_0 \;\Rightarrow\; p_F \sim 500\,\mathrm{MeV} \;\Rightarrow\; f_\pi \sim 100\,\mathrm{MeV}$$

Matching, part II

- Compute m dependence of vacuum energy

$$\Delta\mathcal{E} = -B_1\left[\mathrm{Tr}(M)\right]^2 - B_2\mathrm{Tr}(M^2) \qquad \Sigma = 1$$

- microscopic theory



$$\Delta\mathcal{E} \;\sim\; \left(\frac{g^2}{p_F^4}\right)\left[p_F^2\Delta\log(\Delta)\right]^2\left\{[\mathrm{Tr}(M)]^2 - \mathrm{Tr}(M^2)\right\}$$
$$\sim \Delta^2\left\{[\mathrm{Tr}(M)]^2 - \mathrm{Tr}(M^2)\right\}$$

$$B_1 = -B_2 = \frac{3\Delta^2}{4\pi^2}$$

- meson masses

$$m_\pi^2 \;=\; \frac{3\Delta^2}{4f_\pi^2}(m_u+m_d)m_s$$
$$m_{K^\pm}^2 \;=\; \frac{3\Delta^2}{4f_\pi^2}(m_u+m_s)m_d$$

- Note:

$$m_{GB} \sim 10\ \mathrm{MeV} \qquad m_K < m_\pi$$

Matching, part III

- consider $1/p_F$ expansion

$$\mathcal{L} = \psi_L^\dagger \left(p_0 - \epsilon_p - \frac{MM^\dagger}{2p_F} \right) \psi_L + \frac{\Delta}{2} \psi_L C \psi_L$$

$$+ (L \leftrightarrow R, M \leftrightarrow M^\dagger) + O(1/p_F^2)$$

- $MM^\dagger$ and $M^\dagger M$ enter as gauge fields

$$W_L = \frac{MM^\dagger}{2p_F} \qquad \psi_L \to L\psi_L, \;\; \psi_R \to R\psi_R$$

$$W_R = \frac{M^\dagger M}{2p_F} \qquad W_L \to L W_L L^\dagger + iL\partial_0 L^\dagger, \ldots$$

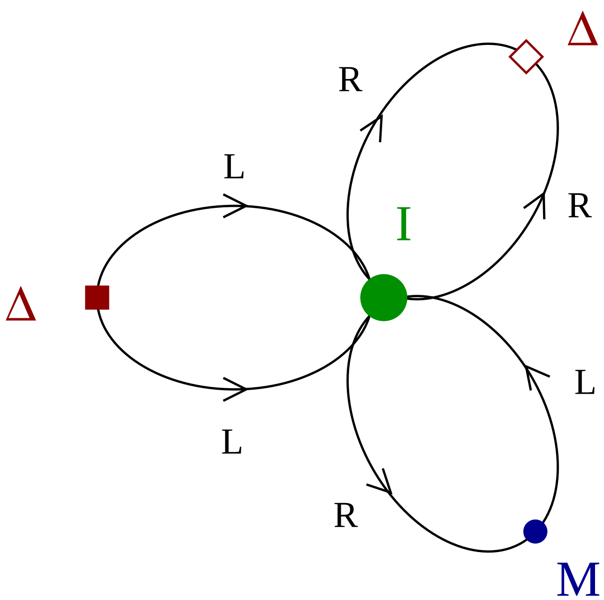
- implement gauge symmetry in effective lagrangian

$$\mathcal{L} = \frac{f_\pi^2}{4} \mathrm{Tr} \left(\nabla_0 \Sigma \nabla_0 \Sigma^\dagger \right) + \ldots$$

$$\nabla_0 \Sigma = \partial_0 \Sigma + i \frac{MM^\dagger}{2p_F} \Sigma - i \Sigma \frac{M^\dagger M}{2p_F}$$

Matching, anomalous part

- linear term $\mathrm{Tr}(M\Sigma)$ in vacuum energy related to instantons (violates $(Z_2)_A$ symmetry)



Instanton size

$$\rho \sim \mu^{-1} \ll \Lambda_{QCD}^{-1}$$

- instanton contribution to vacuum energy: $\mathcal{E} = A\mathrm{Tr}(M)$

$$A = C_{N_c}^{N_f} \langle \psi C \psi \rangle^2 \left(b \log\left(\frac{\mu}{\Lambda_{QCD}}\right)\right)^6 \left(\frac{\Lambda_{QCD}}{\mu}\right)^1 2\Lambda_{QCD}^{-3} e^{i\Theta}$$

$$\langle \psi C \psi \rangle = \frac{3\sqrt{2}\pi}{g} \Delta \left(\frac{\mu^2}{2\pi^2}\right)$$

- Note: $\langle \bar\psi \psi \rangle = -A$

$$\langle \bar\psi \psi \rangle \sim \left(\frac{\Lambda_{QCD}}{\mu}\right)^8 \Lambda_{QCD}^{-3} \ll \Lambda^{-3}$$

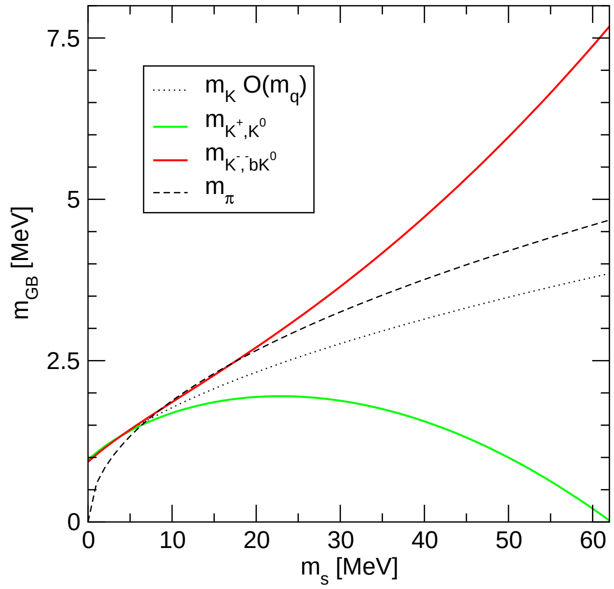
Kaon condensation

- K^0 mass vanishes if

$$m_s > m_s|_{crit}$$

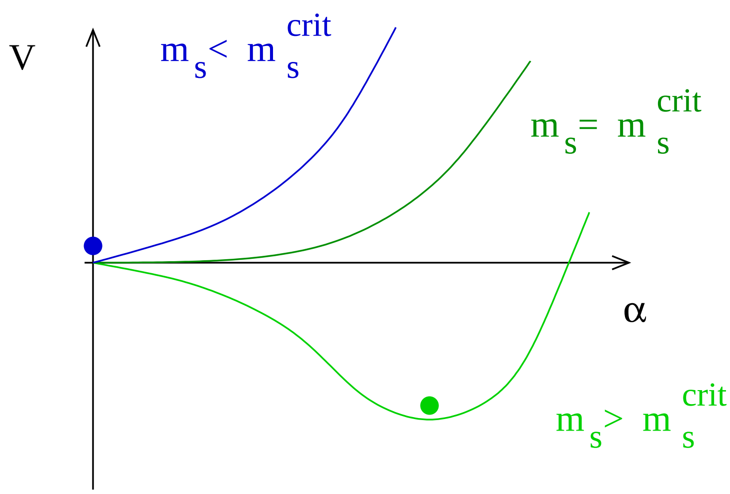
$$m_s|_{crit} \simeq 3 m_u^{1/3} \Delta^{2/3}$$

$\Rightarrow$ kaon condensation



- new vacuum state: $\Sigma = \exp(i\alpha\lambda_4)$

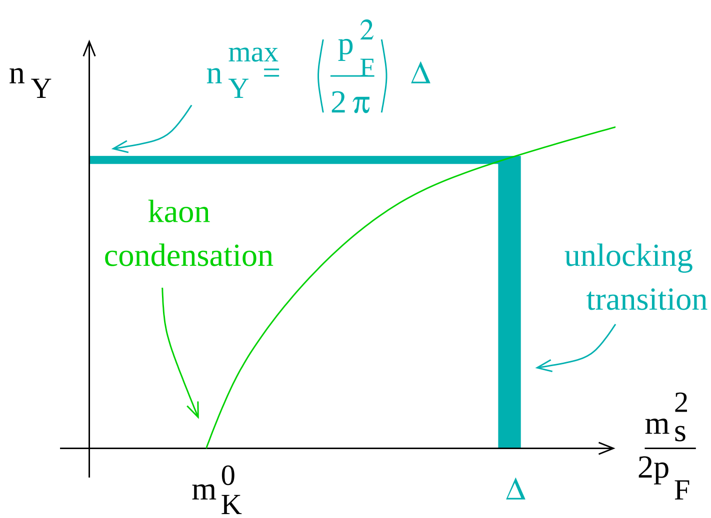
$$V(\alpha) = -f_\pi^2 \{ \frac{m_s^2}{4p_F} \sin^2(\alpha) + (m_K^{(0)})^2 (\cos(\alpha) - 1) \}$$



- hypercharge density

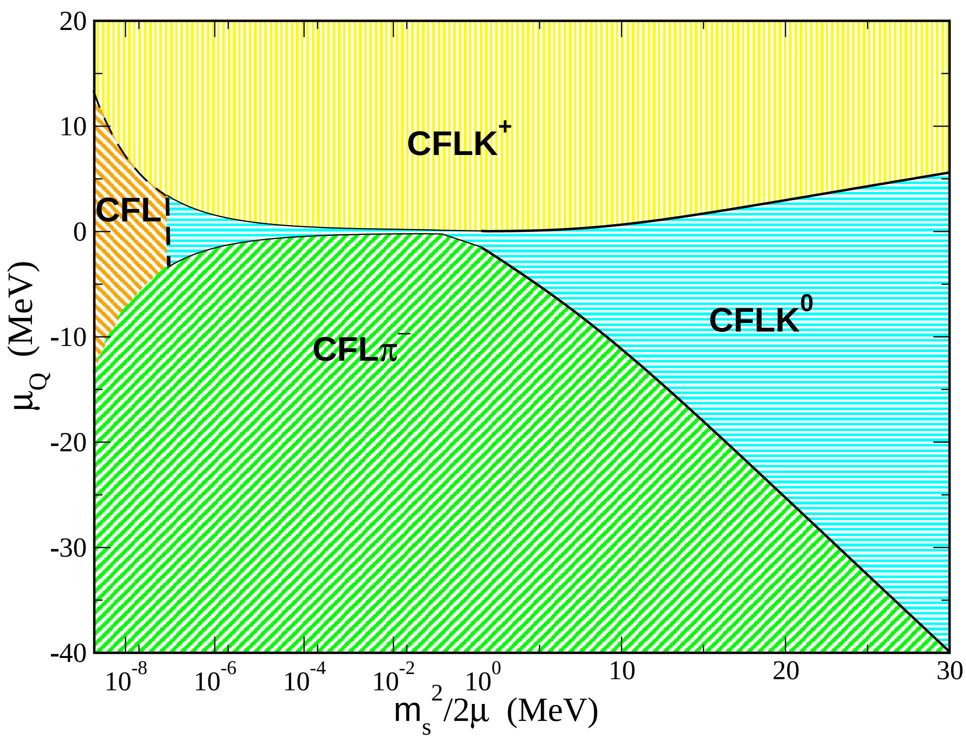
$$n_Y = f_\pi^2 \mu_{eff} \{ 1 - \frac{(m_K^{(0)})^4}{\mu_{eff}^4} \}$$

$$\mu_{eff} = m_s^2 / (2p_F)$$



Phase Diagram of CFL phase

$$V(\Sigma) = \frac{f_\pi^2}{2}\mathrm{Tr}\left(W_L\Sigma W_R\Sigma^\dagger\right) - A\mathrm{Tr}(M\Sigma^\dagger) - B_1\left[\mathrm{Tr}(M\Sigma^\dagger)\right]^2 + \ldots$$



- response to masses/ext. fields determined by collective states

T. Schäfer, P. Bedaque, Nucl. Phys. A697,802,2002

D. Kaplan, S. Reddy, Phys.Rev.D65:054042,2002

- Goldstone modes dominate transport, neutrino emissivity, ...

Jaikumar, Prakash, Schäfer (2002)

Ellis, Shovkovy (2002)

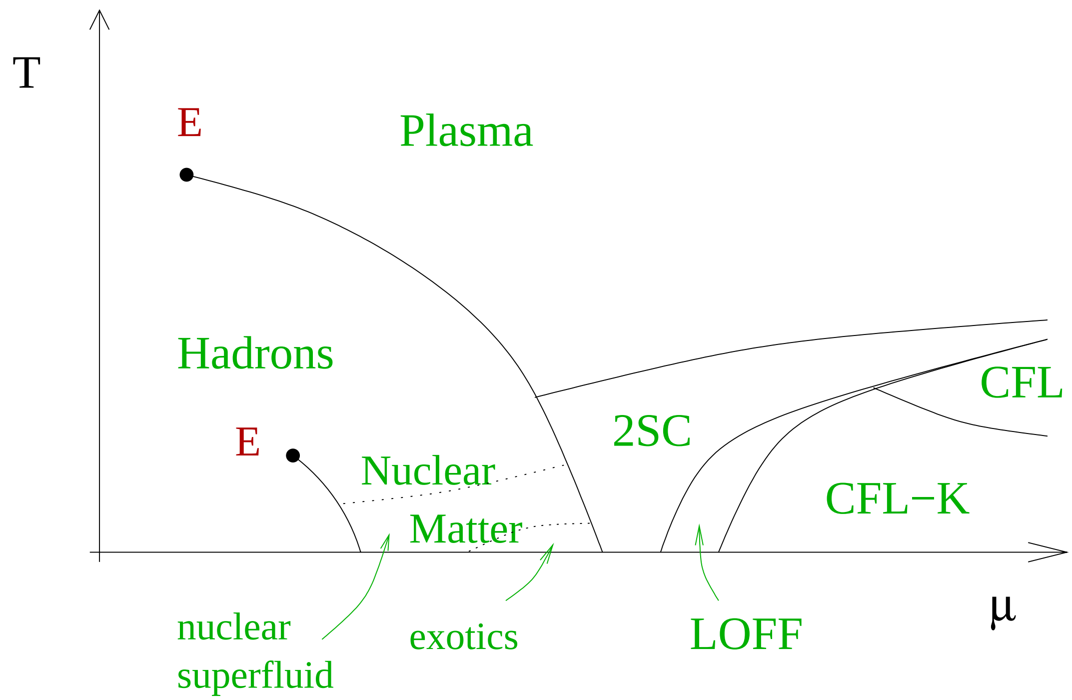
Reddy, Sadzikowski, Tachibana (2002)

- superconducting K-strings

Kaplan, Reddy (2002)

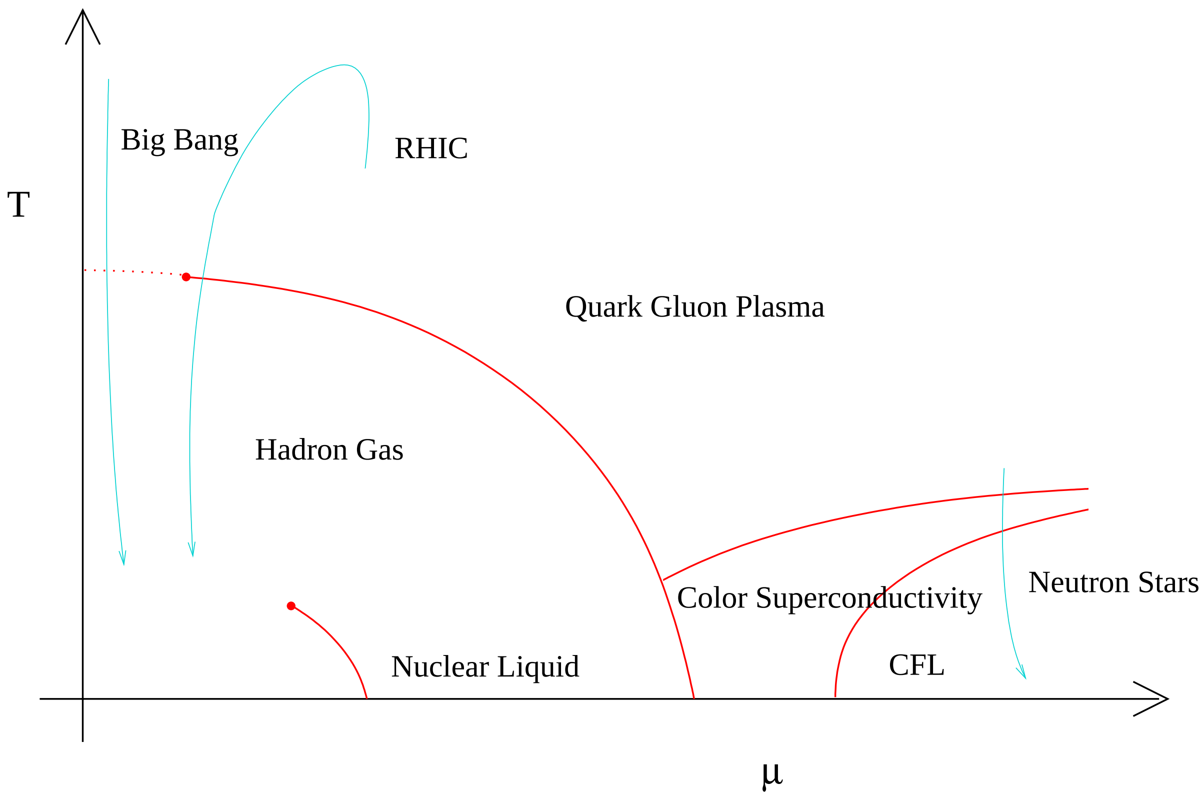
Buckley, Zhitnitsky (2002)

Phase Diagram: $m_s \neq 0$



- phase structure at moderate μ (and $m_s \neq 0$) complicated and poorly understood
- "real world" effects (charge neutrality, neutrino transparacy) important
- observables for neutron star phenomenology

Summary: The Many Phases (Faces?) of QCD



- Interesting new phases of QCD at non-zero baryon density: Color Superconductivity, Color-Flavor-Locking
- Theory: Weak coupling realization of chiral symmetry breaking and mass gaps
- Phenomenology: Neutron Stars (cooling, neutrino burst, glitches?), Heavy Ions (Tri-critical point, precursor phenomena)